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Honeycomb Layer of Cobalt(II) Azide Hydrazine Showing Weak Ferromagnetism

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Using hydrazine (N₂H₄) as a cobridge with azide, a honeycomblayered cobalt(II) coordination polymer, $Co(N_2H_4)(N_3)_2$ (1), is obtained that exhibits spin-canted weak ferromagnetism with T_N of 13.5 K.

Molecule-based magnets bridged by conjugated ligands have experienced continuous interest because of their fascinating physical properties.¹ Also, up to now, a large number of azide (N_3^-)-bridged magnetic compounds have been synthesized and characterized because of not only the extreme versatility of the bridging modes [end-to-end (EE), end-on (EO), and other more complicated modes] but also the exceptional ability to transmit different types of magnetic

coupling.² However, not so many cobalt(II) azide complexes have been reported.³ probably because the oxidation of Co^{II} ions is easy in the presence of N_3^- ions,⁴ especially when alkalescent ligands are employed. A possible approach to avoiding the oxidation of Co^{II} is to supply a reduction environment during the synthesis process. Herein, using hydrazine sulfate (N_2H_4 · H_2SO_4) as the reducing protective agent,5 we prepared a novel 2D honeycomb-like azidebridged cobalt(II) compound, $Co(N_2H_4)(N_3)_2$ (1), in which N₂H₄, in fact, acts as the second bridging ligand through a bis-monodentate coordination mode.⁶ Magnetic studies show that 1 is a weak ferromagnet due to spin-canting, with the critical temperature $T_{\rm N}$ being 13.5 K. To the best of our knowledge, compound 1 is the first example of an azidebridged magnetic polymer that contains short bridging hydrazine as the coligand.

The reaction of $CoSO_4$ ·7H₂O, N₂H₄·H₂SO₄, and NaN₃ in a 1:1:20 molar ratio in water yielded dark-red column crystals of 1.⁷ Single-crystal X-ray analysis⁸ revealed that 1 consists

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⁽⁷⁾ Preparation: All starting chemicals were commercially available reagents and were used without further purification. *Caution!* Azido compounds of metal ions are potentially explosive, and only a small amount of the materials should be prepared and handled with care. An aqueous solution (10 mL) of hydrazine sulfate (0.13 g, 1.0 mmol) and CoSO₄·7H₂O (0.28 g, 1.0 mmol) was heated at 93 °C for 10 min and then quickly mixed with a hot aqueous solution (15 mL) of excessive NaN₃(1.3 g, 20 mmol). The mixed purple solution was kept at 93 °C for 10 min without disturbance. After slow cooling down to room temperature at 5 °C/h, X-ray-quality dark-red column crystals were obtained. The crystals was filtered and washed with distilled water and ethanol, respectively, and then dried in vacuo. Yield: 52% based on Co. Anal. Calcd for 1: H, 2.30; N, 64.02. Found: H, 2.40; N, 63.69. IR bands (cm⁻¹): 2071 (vs) for stretching of the azide, 3319 for ν_{NH2} of hydrazine, 1577 for δ_{NH2} of hydrazine.



Figure 1. View of the coordination environment of **1** along the *ab* plane and the honeycomb layer bridged by EO/EE azide and hydrazine. Color scheme: Co, purple; N, blue. Hydrogen atoms were omitted for clarity.

of neutral honeycomb 2D networks of stoichiometric Co- $(N_2H_4)(N_3)_2$ (Figure 1).

Up to now, several honeycombed cobalt(II) moleculebased magnets have been reported, such as $Co(N_3)_2(bpg)$. DMSO (2; bpg = $meso-\alpha,\beta$ -bis(4-pyridyl) glycol, DMSO = dimethyl sulfoxide)^{3e} and $Co_3Cl_4(H_2O)_2[Co(Hbbiz)_3]_2$ [3; H_2 bbiz = 2,2'-bis(benzimidazole)];⁹ however, they are all different from 1 with regards to the details of connections between cobalt sites. For 2, the Co^{2+} ions form a honeycomb layer just bridged by N₃⁻ ligands, which turns into a square structure if the two ferromagnetically coupled Co²⁺ ions are viewed as one node. For 3, the 2D framework consists of 12-membered rings of alternating tetrahedral and octahedral Co²⁺ ions, creating an extended honeycomb hexagonal lattice. In our compound 1, the network is along the *ab* plane, made up of Co^{II} ions bridged by N_3^- and N_2H_4 . In the network, the two crystallographically independent Co²⁺ ions, Co1 and Co2, possessing distorted CoN₆ octahedral geometries, are linked by one EO azide and one hydrazine in the bis-monodentate mode along the *a* direction, and they are further connected by two EE azido ligands along the bdirection. In other words, the 2D network can be described as double EE azido-linked Co-N₃/N₂H₄ chains, affording the honeycomb layer. The two coordination hydrazine ligands are in trans positions around Co1 but in cis sites around Co2. The Co-N distances are Co1-N = 2.120-2.163 Å and Co2-N = 2.118-2.189 Å. In the $Co-N_3/N_2H_4$ chain, the $Co1-N_{EO azido}-Co2$ angle is 116.6° and the torsion angle Co1-(N-N)_{hydrazine}-Co2 is 50.0°, with an intrachain



Figure 2. Temperature dependence of $\chi_M T$ and χ_M^{-1} of **1** at H = 1 kOe from 2 to 300 K. Inset: ZFC and FC magnetization at H = 20 Oe from 2 to 35 K. The pink line represents the best fit to the Curie–Weiss law, and the blue line corresponds to the best fit with Curély's model.



Figure 3. Hysteresis loop at T = 2 K for **1**. Inset: hysteresis from -1.5 to +1.5 kOe for clarity and the *M* vs *H* plot.

Co···Co separation of 3.606 Å. The interchain Co···Co separation spanned by a double EE azido bridge is 5.173 Å. Therefore, the hexagonal window is elongated with two longer opposite sides.

The 3D structure of **1** is the stacking of these 2D networks or sheets along the *c* direction (Figure S1 of the Supporting Information), and the sheets are glued by the N–H_{hydrazine}• ••N_{EO azido} hydrogen bonds between the sheets, with the hydrogen-bonding geometry of N•••N distance = 3.040 Å and N–H•••N angle = 112.74°. The shortest intersheet Co• ••Co distance is 6.099 Å. It is noted that **1** belongs to a noncentrosymmetric space group of C222₁. There is no inversion center between the neighboring Co²⁺ ions, and the neighboring coordination octahedra are alternately tilted toward each other (Figure S2 of the Supporting Information). These noncentered structural features suggest the possibility of spin-canting by the Dzyaloshinsky–Moriya (D–M) or antisymmetric interaction.¹⁰

Indeed, magnetic studies reveal that the material is a weak ferromagnet at low temperature (Figures 2 and 3 and Figures S3 and S4 of the Supporting Information). Under a 1 kOe field, the $\chi_M T$ value of 3.45 cm³ mol⁻¹ K at 300 K is typical for high-spin Co²⁺ ions with significant orbital contribution. Upon a lowering of the temperature, the $\chi_M T$ values decrease gradually, reach a minimum of 1.52 cm³ mol⁻¹ K at 28 K, and then increase quickly to a maximum of 2.05 cm³ mol⁻¹

⁽⁸⁾ Characterization: The crystallographic data for the single crystal of 1 were collected at 293 K on a Nonius Kappa CCD diffractometer using graphite-monochromated Mo Kα radiation of λ = 0.710 73 Å. The structure was solved by direct methods through the *SHELX-97* programs. Crystal data for 1: Co₂N₁₆H₈, M_w = 350.04, orthorhombic, red column (0.1 × 0.4 × 0.2 mm), space group C2221, a = 6.6121(5) Å, b = 13.2285(10) Å, c = 12.0851(11) Å, V = 1057.06(15) Å³, Z = 4, ρ_{calcd} = 2.199 g/cm³, GOF = 0.923, R1 = 0.0423, wR2 = 0.0694. dc and ac magnetization measurements were performed on a Quantum Design MPMS XL-5 SQUID system. Diamagnetic corrections were estimated using Pascal constants (-87.5 × 10⁻⁶ cm³/mol) and background corrections by experimental measurement on the sample holder.

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K at 14 K. Upon further cooling, $\chi_M T$ drops sharply. This behavior indicates the occurrence of weak ferromagnetism (WF) below ca. 14 K. The susceptibility data above 130 K obeys the Curie–Weiss law, with a Curie constant C = 3.74cm³ mol⁻¹ K and a Weiss constant $\theta = -22.7$ K. The negative θ suggests an overall antiferromagnetic (AF) interaction between the Co²⁺ ions and/or the spin–orbit coupling effect of Co^{II}.

The material was further characterized by zero-field-cooled (ZFC) and field-cooled (FC) measurements, alternating current (ac) susceptibility, and isothermal magnetization in a low-temperature region, and the results confirmed the WF. As shown in the inset in Figure 2, in a field of 20 Oe, ZFC/ FC measurements revealed spontaneous weak magnetization due to the onset of 3D long-range ordering of WF, and the bifurcation of the ZFC and FC plots below the Néel temperature of 13.5 K indicates irreversibility. In ac data (Figure S3 of the Supporting Information), the almost nonfrequency-dependent peaks of χ_{M}' and χ_{M}'' , being at 13.5 and 13.0 K, respectively, are agreeable to the previous observation in direct current (dc) measurement. Finally, the isothermal magnetization at 2 K (Figure 3) first shows a relatively quick initial increase and not a linear increase with the field up to 10 kOe and then goes almost linearly and more flatly up to 0.30 N β per Co²⁺ at 50 kOe, far from the expected saturation value of the Co²⁺ ion, which suggests again the overall AF coupling in the material. The M(H)curve is of somewhat an S shape, which indicates the slight AF interaction between the layers. A hysteresis loop is observed clearly at 2 K with a coercive field $H_c \approx 800$ Oe and a remnant magnetization of $M_r = 0.012 \text{ N}\beta$. According to the above magnetic study, we can conclude that 1 is a weak ferromagnet.

There are two different magnetic interactions within the network (Figure 1), namely, J_1 and J_2 , representing the couplings through the EO-N₃/N₂H₄ bridge and the double EE-N₃ bridge, respectively. Because there is no appropriate

expression for this complicated system, here, we attempt to evaluate the magnitude of the exchange couplings by a model applied for the honeycomb lattice by Curély et al.¹¹ (see the Supporting Information). The values of J_1 and J_2 estimated by this model, +2.16 and -29.9 cm⁻¹, indicate a ferromagnetic coupling through an EO-N₃/N₂H₄ bridge and an AF coupling through double EE-N₃ connections, and the latter is consistent with the values through the same bridges reported in the literature.¹² However, the limitation should be noted that this classical model, viewing J and g as isotropic, is quite approximate because of the great anisotropy for Co^{II}. A more involved theoretical analysis is needed.

As is well-known, the spin-canting is due to the singleion magnetic anisotropy and/or antisymmetric interactions.¹⁰ In the case of **1**, the lack of an inversion center between Co^{2+} sites, connected by EO-N₃/N₂H₄ bridges or double EE-N₃ ones, allows the occurrence of antisymmetric or D–M interaction, and the single-ion anisotropy of Co^{2+} enhances this interaction because the antisymmetric exchange is proportional to the single-ion anisotropy. Therefore, the magnetic behavior observed in **1** is reasonable. The combination of single-ion and structural anisotropies competes with the AF overall ordering and results in canted WF.¹³

In summary, a novel cobalt(II) azide hydrazine complex has been obtained. This is the first example of an azidobridged molecule-based magnet with N₂H₄ as the coligand. In the structure, the Co²⁺ ions are connected by EO-N₃/N₂H₄ or double EE-N₃ bridges into a 2D honeycomb network. The material exhibits WF behavior with T_N of 13.5 K due to the absence of an inversion center between Co²⁺ sites and the large anisotropy of the Co²⁺ ions. Further studies will explore analogues of other metals and their magnetism, and these are in progress.

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Supporting Information Available: X-ray crystallographic file in CIF format for **1**, Figures S1–S4, and the model of the honeycomb lattice. This material is available free of charge via the Internet at http://pubs.acs.org.

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